Boundedness of Hardy-Steklov operator on weighted Lorentz Spaces $\Lambda_u^p(w)$

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Abstract.

In this paper, we obtain the necessary and sufficient conditions on pair of weights (w_0, w_1) for the boundedness of the Hardy-Steklov operator $\int_{a(x)}^{b(x)} f(t) dt$ between the spaces $\Lambda_u^p(w_0)$ and $\Lambda_v^q(w_1)$ for non-negative functions f, where 1 .

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1. Introduction

By a weight function u on $(0,\infty)$ we mean a non-negative locally integrable measurable function. We take $M_0^+ \equiv M_0^+((0,\infty), u(x)dx)$ to be the set of all those functions which are measurable, non-negative and finite a.e. on $(0,\infty)$ with respect to the measure u(x)dx. Then the *distribution function* f_{*_u} of $f \in M_0^+$ is given by

$$f_{*_{u}}(\lambda) = \int_{\{x \in (0,\infty): f(x) > \lambda\}} u(x) dx, \quad \lambda \ge 0.$$

For a measurable set E, by u(E) we mean $\int_{E} u(x) dx$.

The non-increasing rearrangement f_u^* of f is defined as

$$f_u^*(t) = \inf \{ \lambda : f_{*_u}(\lambda) \le t \}, \quad t \ge 0.$$

Essentially, $f_{*_{u}}^{-1} = f_{u}^{*}$ [3].

The weighted Lorentz spaces $\Lambda_{v}^{q}(w)$ are defined to be the collection of those members of $M_{0}^{+}((0,\infty),v(x)dx)$ for which

$$||f||_{\mathbf{A}^{q}_{\mathcal{V}}(w)} = \begin{cases} \int_{0}^{\infty} f_{\nu}^{*q}(x)w(x)dx, & 0 < q < \infty \\\\ \sup_{t>0} \left(\int_{0}^{t} w(x)dx \right)^{1/q} f_{\nu}^{*}(t), & q = \infty \end{cases}$$

is finite.

For $w(x) = \frac{q}{p} x^{\frac{q}{p}-1}$, the spaces we get are called *two exponent Lorentz spaces*, denoted as $L_{\nu}^{p,q}(0,\infty)$, consisting of $f \in \mathsf{M}_{0}^{+}$ for which

$$\|f\|_{L^{p,q}_{v}(0,\infty)} = \begin{cases} \int_{0}^{\infty} \frac{p}{q} [t^{1/p} f_{v}^{*}(t)]^{q} \frac{dt}{t}, & 0 < q < \infty \\\\ \sup_{t > 0} t^{1/p} f_{v}^{*}(t), & q = \infty \end{cases}$$

is finite. For p = q, the spaces are reduced to *Lebesgue spaces* $L_{\nu}^{p}(0,\infty)$. Again, for $\nu = 1$, the spaces $\Lambda_{\nu}^{q}(w)$ become $\Lambda^{q}(w)$, known as *classical Lorentz spaces*. Further, on taking $w(x) = x^{\frac{q}{p}-1}$, we get the so called *two exponent classical Lorentz spaces* $L^{p,q}$. The spaces $\Lambda^{q}(w)$, may also be looked upon as $L^{q}(w)$ for the functions f^{*} *i.e.*, $\|f\|_{\Lambda^{q}(w)} = \|f^{*}\|_{L^{q}(w)}$. And, in case, f is non-increasing, it is precisely $L^{q}(w)$. In fact, each of the Lorentz spaces talked here, may be deduced from $\Lambda_{\nu}^{q}(w)$, $0 < q \le \infty$. The spaces $\Lambda_{\nu}^{q}(w)$ have been considered for studying the boundedness of various operators. E.g., the boundedness of Hardy Littlewood maximal function on $\Lambda_{\nu}^{q}(w)$ by Carro and Soria [2] etc.

In the present paper, we study the boundedness of Hardy-Steklov operator on weighted Lorentz spaces. In Section 2, we give some of the known results (from [2], [4] and [6]) which will be used in the subsequent Section 3, containing the main result.

2. Known Results

Definition 1. The Hardy-Steklov operator is defined as

$$(Tf)(x) = \int_{a(x)}^{b(x)} f(t)dt$$

where the functions a = a(x) and b = b(x) are strictly increasing and differentiable on $(0, \infty)$. Also, they satisfy

 $a(0) = b(0) = 0; \quad a(\infty) = b(\infty) = \infty \text{ and } a(x) < b(x) \text{ for } 0 < x < \infty.$

Clearly a^{-1} and b^{-1} exist, and are strictly increasing and differentiable as well, for the functions introduced in the definition of Hardy-Steklov operator.

Definition 2. A function f is said to satisfy Δ_2 -condition if $f(2x) \le Cf(x) \forall x > 0$, for some constant C > 0.

Theorem A. $\|\cdot\|_{\mathbf{A}^q_{\mathcal{V}}(w)}$ is a quasi-norm if and only if $\int_0^x w(t) dt$ satisfies Δ_2 -condition.

Notations.

$$W(y) = \int_{a(x)}^{y} w(t)dt, \quad \widetilde{W}(y) = \int_{y}^{b(x)} w(t)dt;$$
$$u^{a}(y) = u(a^{-1}(y))(a^{-1})'(y), \quad u^{b}(y) = u(b^{-1}(y))(b^{-1})'(y).$$

We consider now, a sequence $\{m_k\}_{k \in \mathbb{Z}}$ defined as follows: Let m > 0 be fix, set $m_0 = m$, and

$$m_{k+1} = a^{-1}(b(m_k)), \quad k \ge 0$$

$$m_k = b^{-1}(a(m_{k+1})), \quad k < 0.$$
(1)

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Obviously $a(m_{k+1}) = b(m_k) \quad \forall \quad k \in \mathsf{Z}$.

Lemma 1. Fix m > 0 and define $\{m_k\}_{k \in \mathbb{Z}}$ as above. Then

$$m_k < m_{k+1} \text{ for } k \in \mathsf{Z}; \quad \lim_{k \to \infty} m_k = \infty; \text{ and } \lim_{k \to -\infty} m_k = 0$$

Lemma 2. Let $p \le q$. Then a weight w_0 satisfies that for every $\{t_k\}_k \subset \mathsf{R}^+$,

$$\left(\sum_{k} \left(\int_{0}^{t_{k}} w_{0}(s) ds\right)^{q/p}\right)^{p/q} \leq C \int_{0}^{\sum t_{k}} w_{0}(s) ds,$$

$$(2)$$

if and only if, for every collection of functions $\{f_k\}_k$ in $\Lambda^p_u(w_0)$ with pairwise disjoint support, there exists a constant C > 0 such that

$$\sum_{k} \left\| f_{k} \right\|_{\mathbf{\Lambda}^{p}_{u}(w_{0})}^{q} \leq C \left\| \sum_{k} f_{k} \right\|_{\mathbf{\Lambda}^{p}_{u}(w_{0})}^{q}.$$

A trivial choice of weights for which (2) holds is to take w_0 to be a non-decreasing weight [2].

Theorem B. (Sawyer's Duality Principle) Let 1 , <math>v(x) and g(x) are nonnegative functions on $[0, \infty[$ with v locally integrable. Then

$$\sup_{\substack{\downarrow f \ge 0}} \frac{\int_0^\infty f(x)g(x)dx}{\left(\int_0^\infty f^p(x)v(x)dx\right)^{1/p}} \approx \left(\int_0^\infty \left(\int_0^\infty g\right)^{p'} \left(\int_0^x v\right)^{-p'}v(x)dx\right)^{1/p'} + \left(\int_0^\infty g\right) \left(\int_0^\infty v\right)^{-1/p}$$

where the symbol \approx means that the ratio of left and right hand sides is bounded between two positive constants depending only on p (and not on v or g).

3. Main Result

Theorem. Let 1 , <math>x < a(x) for $0 < x < \infty$, $W_1 \in \Delta_2$ and w_0 be satisfying the condition (2). Then $T : \Lambda^p_u(w_0) \to \Lambda^q_v(w_1)$, i.e., the inequality

$$\left(\int_{0}^{\infty} \left(\int_{a(x)}^{b(x)} f(t)dt\right)_{v}^{*q} w_{1}(x)dx\right)^{1/q} \le C \left(\int_{0}^{\infty} f_{u}^{*p}(x)w_{0}(x)dx\right)^{1/p}$$
(3)

holds if and only if $\forall f \ge 0$

$$\sup_{a(x) < y < b(x)} \left[\int_{a(x)}^{b(x)} \left(\int_{a(x)}^{y} (\chi_{(a(x), y)} u^{-1})_{u}^{*}(t) dt \right)^{p'} W_{0}(y)^{-p'} w_{0}(y) dy + (y - a(x)) W_{0}(b(x))^{1/p} \left[W_{1}(\Phi(y)) \right]^{1/q} < \infty;$$
(4)

where

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$$\Phi(y) = \begin{cases} U_1(y) & \text{when } W_1(z) = \int_{a(x)}^z w_1^b(s) ds \\ \\ \widetilde{U}_1(y) & \text{when } W_1(z) = \int_{a(x)}^z w_1^a(s) ds. \end{cases}$$

Proof. (\Leftarrow) (Sufficiency part)

Fix m > 0 and define $\{m_k\}_{k \in \mathbb{Z}}$ as in (1). Denote $E_k = (m_k, m_{k+1})$, and $a_k = a(m_k), b_k = b(m_k)$. Then for $x \in E_k$, $a(x) < a_{k+1} = b_k < b(x)$. Now since W_1 satisfies Δ_2 -condition, we have

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$$\left(\int_{0}^{\infty} \left(\int_{a(x)}^{b(x)} f(t)dt\right)_{v}^{*q} w_{1}(x)dx\right)^{1/q} = \left(\int_{0}^{\infty} \left(\sum_{k\in\mathbb{Z}}\chi_{E_{k}}(x)\int_{a(x)}^{b_{k}} f(t)dt + \sum_{k\in\mathbb{Z}}\chi_{E_{k}}(x)\int_{a_{k+1}}^{b(x)} f(t)dt\right)_{v}^{*q} w_{1}(x)dx\right)^{1/q} \\ \leq C \left[\underbrace{\left(\int_{0}^{\infty} \left(\sum_{k\in\mathbb{Z}}\chi_{E_{k}}(x)\int_{a(x)}^{b_{k}} f(t)dt\right)_{u}^{*q} w_{1}(x)dx\right)^{1/q}}_{I_{2}} + \underbrace{\left(\int_{0}^{\infty} \left(\sum_{k\in\mathbb{Z}}\chi_{E_{k}}(x)\int_{a_{k+1}}^{b(x)} f(t)dt\right)_{v}^{*q} w_{1}(x)dx\right)^{1/q}}_{I_{1}}\right]. \quad (5)$$

Now notice that, for $x = m_{k+1}$ the condition (4) becomes

$$\sup_{a_{k+1} < y < b_{k+1}} \left[\int_{a_{k+1}}^{b_{k+1}} \left(\int_{a_{k+1}}^{y} (\chi_{(a_{k+1}, y)} u^{-1})_{u}^{*}(t) dt \right)^{p'} \left(\int_{a_{k+1}}^{y} w_{0}(t) dt \right)^{-p'} w_{0}(y) dy + (y - a_{k+1}) \left(\int_{a_{k+1}}^{b_{k+1}} w_{0}(t) dt \right)^{1/p} \right] \left(\int_{a_{k+1}}^{y} u_{1}(t) dt w_{1}^{b}(s) ds \right)^{1/q} < \infty$$
(6)

which is a sufficient condition for the inequality

$$\left(\int_{a_{k+1}}^{b_{k+1}} \left(\int_{a_{k+1}}^{y} f(t)dt\right)_{v}^{*q} w_{1}^{b}(y)dy\right)^{1/q} \le C \left(\int_{a_{k+1}}^{b_{k+1}} f_{u}^{*p}(x)w_{0}(x)dx\right)^{1/p}$$
(7)

to hold. Thus, in view of (6), (7) and Lemma 2, we have

$$I_{1} = \left(\sum_{k \in \mathbb{Z}} \int_{m_{k}}^{m_{k}+1} \left(\int_{a_{k+1}}^{b(x)} f(t)dt\right)_{v}^{*q} w_{1}(x)dx\right)^{1/q}$$

$$= \left(\sum_{k \in \mathbb{Z}} \int_{a_{k+1}}^{b_{k+1}} \left(\int_{a_{k+1}}^{y} f(t)dt\right)_{v}^{*q} w_{1}^{b}(y)dy\right)^{1/q}$$

$$\leq C \left(\sum_{k \in \mathbb{Z}} \left(\int_{a_{k+1}}^{b_{k+1}} f_{u}^{*p}(x)w_{0}(x)dx\right)^{q/p}\right)^{1/q}$$

$$= C \left(\sum_{k \in \mathbb{Z}} \left(\int_{0}^{\infty} (f_{k})_{u}^{*p}(x)w_{0}(x)dx\right)^{q/p}\right)^{1/q} \text{ for } f_{k} = \chi_{(a_{k+1},b_{k+1})}f$$

$$\leq C \left(\int_{0}^{\infty} (\sum_{k \in \mathbb{Z}} f_{k})_{u}^{*p}(x)w_{0}(x)dx\right)^{1/p}$$

$$= C \left(\int_{0}^{\infty} f_{u}^{*p}(x)w_{0}(x)dx\right)^{1/p}.$$
(8)

Again, on taking $x = m_k$ in (4), we obtain the expressions conjugate to (6) and (7). And then, further, using Lemma 2, we get

$$I_{2} = \left(\sum_{k \in \mathbb{Z}} \int_{m_{k}}^{m_{k}+1} \left(\int_{a(x)}^{b_{k}} f(t) dt \right)_{v}^{*q} w_{1}(x) dx \right)^{1/q}$$

$$= \left(\sum_{k \in \mathbb{Z}} \int_{a_{k}}^{b_{k}} \left(\int_{y}^{b_{k}} f(t) dt \right)_{v}^{*q} w_{1}^{b}(y) dy \right)^{1/q}$$

$$\leq C \left(\sum_{k \in \mathbb{Z}} \left(\int_{a_{k}}^{b_{k}} f_{u}^{*p}(x) w_{0}(x) dx \right)^{q/p} \right)^{1/q}$$

$$\leq C \left(\int_{0}^{\infty} f_{u}^{*p}(x) w_{0}(x) dx \right)^{1/p}.$$
(9)

The inequality (3), now follows from (8) and (9).

 (\Rightarrow) (Necessary part)

It can be easily verified that if the condition (2) is true for a weight w_0 , then so is this for $\tilde{w}_0 = w_0 \chi_{(a(x),b(x))}$, $x \in (0,\infty)$. So, let $f \in \Lambda^p_u(\tilde{w}_0)$, Then, for every $t \in [y, b(x)]$, we have

$$\int_{a(x)}^{y} f(s) ds \leq T(\chi_{(a(x),y)} f)(t).$$

Let $\eta < \int_{a(x)}^{y} f(s) ds$, then

$$\int_{y}^{b(x)} v(s) ds \leq [T(\chi_{(a(x),y)}f)]_{*_{v}}(\eta).$$

Thus,

$$\begin{split} \eta \bigg(\int_{a(x)}^{\int_{y}^{b(x)} v(s) ds} w_{1}^{a}(s) ds \bigg)^{1/q} &= \eta \bigg(\int_{x}^{a^{-1} (\int_{y}^{b(x)} v(s) ds)} w_{1}(z) dz \bigg)^{1/q} \\ &\leq \eta \bigg(\int_{x}^{\int_{y}^{b(x)} v(s) ds} w_{1}(z) dz \bigg)^{1/q} \\ &\leq \eta \bigg(\int_{x}^{[T(\chi(a(x), y)f)]_{*_{v}}(\eta)} w_{1}(z) dz \bigg)^{1/q} \\ &\leq \sup_{z>0} z \bigg(\int_{[T(\chi(a(x), y)f)]_{*_{v}}(z)} w_{1}(s) ds \bigg)^{1/q} \\ &= \big\| T(f\chi_{(a(x), y)}) \big\|_{\mathbf{A}_{v}^{q;\infty}(w_{1})} \\ &\leq \big\| T(f\chi_{(a(x), y)}) \big\|_{\mathbf{A}_{v}^{q;\infty}(w_{1})} \\ &\leq C \big\| f \big\|_{\mathbf{A}_{u}^{p}(\widetilde{w}_{0})}. \end{split}$$

Now, on taking the supremum over all $\eta < \int_{a(x)}^{y} f(s) ds$, and using the Sawyer's duality principle, we obtain the condition (4) with $W_1(z) = \int_{a(x)}^{z} w_1^a(s) ds$.

Again for every $t \in [a(x), y]$, we have

$$\int_{y}^{b(x)} f(s) ds \leq T(\chi_{(y,b(x))}f)(t),$$

and therefore, on taking $\eta < \int_{y}^{b(x)} f(s) ds$, we get

$$\int_{a(x)}^{y} v(s) ds \leq \int_{\{s:T(\chi(y,b(x)),f)(s) > \eta\}} v(s) ds = [T(\chi_{(y,b(x))},f)]_{*_{v}}(\eta).$$

Thus

$$\eta \left(\int_{a(x)}^{\int_{a(x)}^{y} v(s) ds} w_1^b(s) ds \right)^{1/q} \le \eta \left(\int_{x}^{\int_{a(x)}^{y} v(s) ds} w_1(z) dz \right)^{1/q}$$

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$$\leq \sup_{z>0} z \left(\int_0^{[T(\chi(y,b(x))f)]_{*_v}(z)} w_1(s) ds \right)^{1/q} \leq C \|f\|_{\mathbf{\Lambda}^p_u}(\widetilde{w}_0).$$

Taking supremum over all those $\eta < \int_{y}^{b(x)} f(s) ds$, we have

$$\sup_{a(x) < y < b(x)} \left(\sup_{f} \frac{\int_{a(x)}^{y} f(t) dt}{\|f\|_{\mathbf{A}^{p}_{u}(\tilde{w}_{0})}} \right) \left(\int_{a(x)}^{\int_{a(x)}^{y} u_{1}(s) ds} w_{1}^{b}(s) ds \right)^{1/q} < \infty .$$

And, now using Sawyer's duality principle, the condition (4) holds with $W_1(z) = \int_{a(x)}^{z} w_1^b(s) ds$.

Remark. It is known that, in particular, for the two exponent classical Lorentz spaces the following embeddings hold when $p \le q$:

$$L^{p,1} \subset \ldots \subset L^p \subset \ldots \subset L^{p,q} \subset \ldots \subset L^{p,\infty}.$$

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